# Uncertainty Inequality for the Entanglement of Both Two-rebits and Two-qubits States 

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Received: 20 September 2018 / Accepted: 27 June 2019 / Published online: 19 July 2019 © Springer Science+Business Media, LLC, part of Springer Nature 2019


#### Abstract

We assume that both the concurrence $C_{r}$ of a two-rebits state and the concurrence $C_{q}$ of the usual two-qubits states are represented by hermitian operators (observables). We calculate the respective uncertainty $\Delta C_{r}$ and the uncertainty $\Delta C_{q}$ measured both as the standard deviation. We make the strictly mathematical assumption that there exists a canonical conjugate variable (called $\xi$ ) to the concurrence ( $C$ ) such that both quantities satisfy a Robertson's [1] uncertainty inequality of the form $(\Delta A)^{2}(\Delta B)^{2}>\left|\frac{1}{2}\langle[A, B]\rangle\right|^{2}$. From such inequality we impose bounds for both uncertainties $\Delta \xi_{r}$ and $\Delta \xi_{q}$.


Keywords Entanglement • Rebits • Qubits • Uncertainty

Entanglement distinguishes substantially quantum mechanics from classical mechanics [2]. Quantum entanglement plays a key role in superdense coding [3], quantum teleportation [4], and quantum information [5-8]. For a two qubits system (in pure and mixed states), entanglement has been intensively studied in the past [9-13]. Concurrence can be considered as a measure of entanglement of a two-qubits state [9]. In the present work it is studied the uncertainty in the concurrence of both two-rebits and two qubits. In the literature it has not been considered uncertainty for a hypothetical conjugate variable to the concurrence [14, 15].

For quantum mechanics defined over vector real spaces the simplest state is the rebit state which is defined as

$$
\begin{equation*}
|\eta\rangle=a_{0}|0\rangle+a_{1}|1\rangle, \tag{1}
\end{equation*}
$$

where $a_{i}(i=0,1)$ are real numbers such that $a_{0}^{2}+a_{1}^{2}=1$ and $\{|0\rangle,|1\rangle\}$ are the classical bits. It is worth mentioning that for an arbitrary two-rebits system in Ref. [16] it was found an expression for the entanglement of formation. Pure states of two-rebits systems are described by 9 quantities. On the other hand, the usual one-qubit state is defined as

$$
\begin{equation*}
|\psi\rangle=c_{0}|0\rangle+c_{1}|1\rangle, \tag{2}
\end{equation*}
$$

[^0]where $c_{i}(i=0,1)$ are a complex numbers such that $\left|c_{0}\right|^{2}+\left|c_{1}\right|^{2}=1$ and $\{|0\rangle,|1\rangle\}$ are the basic qubits. By the way, the corresponding space of for a two-qubits state is 15 dimensional.

In the present work we calculate the uncertainty of the concurrence $\Delta C$ associated to a pure both two-rebits and two-qubits states. For the above we assume that the concurrence is represented by a hermitian operator (observable). Furthermore, we make a strict mathematical assumption that there exists a hermitian canonical conjugate variable, denoted by $\xi$, such that this satisfies a Robertson's uncertainty relation [1] of the form

$$
\begin{equation*}
\Delta C \Delta \xi \geq \frac{1}{2}|\langle[C, \xi]\rangle|=\frac{1}{2} \tag{3}
\end{equation*}
$$

Let us observe that the above equation might indicate the existence of a hidden 'uncertainty principle' that would involve the concurrence $C$ and a hypothetical complementary quantity $\xi(\hbar=1)$.

For a two qubits system, entanglement is an associated intrinsically quantum quantity. This has been very well studied in the past. It is well known that concurrence of a two-qubit system is a good measure of the entanglement of formation $E_{q}(\rho)$ [9] which is defined as

$$
\begin{equation*}
E_{q}[\rho]=h\left(\frac{1+\sqrt{1-C_{q}^{2}}}{2}\right) \tag{4}
\end{equation*}
$$

wheres

$$
\begin{equation*}
h(x)=-x \log _{2} x-(1-x) \log _{2}(1-x) . \tag{5}
\end{equation*}
$$

The concurrence for a pure two-qubits states is given by

$$
\begin{equation*}
C_{q}=\max \left(0, \lambda_{1}-\lambda_{2}-\lambda_{3}-\lambda_{4}\right), \tag{6}
\end{equation*}
$$

being $\lambda_{i},(i=1, \ldots, 4)$ the square root in decreasing order of the eigenvalues of the matrix $\rho \tilde{\rho}$, with

$$
\begin{equation*}
\tilde{\rho}=\left(\sigma_{y} \otimes \sigma_{y}\right) \rho^{*}\left(\sigma_{y} \otimes \sigma_{y}\right) \tag{7}
\end{equation*}
$$

For a two-qubits system the states become less entangled as the degree of mixtures increases. Here restrict ourselves to a pure two-qubit states.

For a two-rebits system in Ref. [16] it was found a formula for the entanglement of formation of two-rebits state which is given by

$$
\begin{equation*}
E\left[\rho_{r}\right]=h\left(\frac{1+\sqrt{1-C_{r}^{2}}}{2}\right), \tag{8}
\end{equation*}
$$

where

$$
\begin{equation*}
h(x)=-x \log _{2} x-(1-x) \log _{2}(1-x) . \tag{9}
\end{equation*}
$$

The above expressions are identical to the case of two-qubits except that the concurrence of a two-rebits is

$$
\begin{equation*}
C\left[\rho_{r}\right]=\left|\operatorname{tr}\left(\rho_{r} \sigma_{y} \otimes \sigma_{y}\right)\right| . \tag{10}
\end{equation*}
$$

For a two-rebits state the entanglement of formation is the expectation value of one single observable $\sigma_{y} \otimes \sigma_{y}$ A pure two-rebits state can be writen as

$$
\begin{equation*}
\left|\Psi_{r}\right\rangle=\sum_{i=1}^{4} a_{i}\left|\phi_{i}\right\rangle, \tag{11}
\end{equation*}
$$

where we have used decimal notation, $a_{i}$ are real numbers such that

$$
\begin{equation*}
\sum_{i=1}^{4} a_{i}^{2}=1, a_{i} \in \mathfrak{R} . \tag{12}
\end{equation*}
$$

The states $\left|\phi_{i}\right\rangle$ are eigenstates of the operator $\sigma_{y} \otimes \sigma_{y}$. In Ref. [17] it was found an analytical expression for the probability density $P\left(C_{r}^{2}\right)$ of finding a pure two-rebits state with a squared concurrence $C_{r}^{2}$. This has the form

$$
\begin{equation*}
P\left(C_{r}^{2}\right)=\frac{1}{2 \sqrt{C_{r}^{2}}} . \tag{13}
\end{equation*}
$$

The above expression has a strong divergence at $C_{r}^{2}=0$. The domain of (13) is $0 \leq$ $C_{r}^{2} \leq 1$. Concerning to a pure two-qubit state, the distribution of probability has been found analytically in [18] and it is

$$
\begin{equation*}
P\left(C_{q}^{2}\right)=\frac{3}{2} \sqrt{1-C_{q}^{2}} \tag{14}
\end{equation*}
$$

We note that the densities of probability (13) and (14) for two-rebits and two-qubits respectively are even functions of the concurrences. That is, they admit negative values of the concurrences, that is, $P\left(\left(-C_{r}\right)^{2}\right)=P\left(C_{r}^{2}\right)$ and $P\left(\left(-C_{q}\right)^{2}\right)=P\left(C_{q}^{2}\right)$. A negative concurrence has a non trivial interpretation. By the way, the discovery that quantum knowledge can be negative was made by Horodecki-Oppenheim-Winter [19]. They argue that in the quantum world there are things we just cannot know. For instance, we cannot know both the position (energy) and momentum (time) of a microscopic quantum system.

The uncertainty of a variable $X$ is defined as its standard deviation, that is

$$
\begin{equation*}
\Delta X=\sqrt{<X^{2}>-<X>^{2}} \tag{15}
\end{equation*}
$$

where the average value of a function $f(X)$ is

$$
\begin{equation*}
<f(X)>=\int_{0}^{1} f(X) P(X) d X \tag{16}
\end{equation*}
$$

being $P(X)$ the distribution of probability of the variable $X$.

- Uncertainty for the concurrence of a two-rebits states $\left(\Delta C_{r}\right)$

Rebits are a powerful tool for Quantum Information Processing. For instance in Ref. [20] it was concluded that rebits do not affect universality for quantum computation. Rebits may also be defined as a real density matrices of $n$ two-level systems.

By using (13) it is possible to define the expectation value of $C_{r}$ as

$$
\begin{align*}
<C_{r}> & :=\int_{0}^{1} \sqrt{C_{r}^{2}} P\left(C_{r}^{2}\right) d\left(C_{r}^{2}\right) \\
& =\int_{0}^{1} \sqrt{C_{r}^{2}} \frac{1}{2 \sqrt{C_{r}^{2}}} d\left(C_{r}^{2}\right) \\
& =\frac{1}{2} \tag{17}
\end{align*}
$$

The respective expectation value of $C_{r}^{2}$ is

$$
\begin{align*}
<C_{r}^{2}> & =\int_{0}^{1} C_{r}^{2} P\left(C_{r}^{2}\right) d\left(C_{r}^{2}\right) \\
& =\int_{0}^{1} C_{r}^{2} \frac{1}{2 \sqrt{C_{r}^{2}}} d\left(C_{r}^{2}\right) \\
& =\frac{1}{3} \tag{18}
\end{align*}
$$

Substituting (17) and (18) in (15) the uncertainty for the concurrence for a pure two-rebits state is

$$
\begin{align*}
\Delta C_{r} & =\frac{\sqrt{12}}{12} \\
& =0.29 . \tag{19}
\end{align*}
$$

If we assume that the uncertainty $\Delta C_{r}$ satisfies an uncertainty principle such as (3) then the uncertainty of its complementary variable $\xi_{r}$ satisfies

$$
\begin{equation*}
\Delta \xi_{r} \geq \frac{\sqrt{12}}{2}=1.7 \tag{20}
\end{equation*}
$$

To interpret $\xi_{r}$ is a nontrivial task. In fact, this is a challenge for Quantum Information Theory.

- Uncertainty for the concurrence of a two-qubits states $\left(\Delta C_{q}\right)$

By using (14) the expectation value of $C_{q}$ is

$$
\begin{align*}
<C_{q}> & :=\int_{0}^{1} \sqrt{C_{q}^{2}} P\left(C_{q}^{2}\right) d\left(C_{q}^{2}\right) \\
& =\frac{3 \pi}{16} \tag{21}
\end{align*}
$$

The respective expectation value of $C_{q}^{2}$ is

$$
\begin{align*}
<C_{q}^{2}> & =\int_{0}^{1} C_{q}^{2} P\left(C_{q}^{2}\right) d\left(C_{q}^{2}\right) \\
& =\frac{2}{5} \tag{22}
\end{align*}
$$

Substituting (21) and (22) in (15) the uncertainty for the concurrence for a pure two-qubits state is

$$
\begin{align*}
\Delta C_{q} & =\sqrt{\frac{2}{5}-\frac{9 \pi^{2}}{256}} \\
& =0.23 \tag{23}
\end{align*}
$$

If we assume that there exists a hidden uncertainty principle as (3), the uncertainty in the respective complementary variable $\xi_{q}$ is

$$
\begin{equation*}
\Delta \xi_{q} \geq \frac{1}{2 \sqrt{\frac{2}{5}-\frac{9 \pi^{2}}{256}}}=2.2 \tag{24}
\end{equation*}
$$

The interpretation of the variable $\xi_{q}$ would remain also as an open question for Quantum Information Science.

We have calculated the uncertainty of the concurrence associated to a pure two-rebits state and also to a pure two-qubits state. Both concurrences have a non zero value which means that their measure cannot be done with precision. This confirms that they are a quantum mechanics quantities that cannot be known with absolute certainty. We observe that both uncertainties have a similar values i.e. $\Delta C_{r}=0.29$ and $\Delta C_{q}=0.23$. To assume from the point of view strictly mathematical that there exists a complementary quantity $\xi$ to the concurrence $C$ in such a way that both variables satisfy an uncertainty principle (3) is a mere conjecture that requires many efforts and intense search. By choosing a system of units where $\hbar=1$ it is concluded that the complementary variable $\xi$ of (3) is dimensionless. A future challenge is to identify the hermitian operator $\xi$ with a physical quantity.

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